

# Phys 401: Midterm Exam 1

Fall 2017

- Write your name and Student ID number in the space provided below and sign.

Name, Last Name:	
ID Number:	
Signature:	

- You have 2.5 hours.
  - You must show the details of all your work. Illegible and ambiguous explanations and calculations will lead to deduction of your grade.
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**Problem 1** Consider a point particle of mass  $m$  that moves on a straight line under the influence of a conservative force  $F = F(x)$  where  $x$  denotes the position of the particle.

**1.a** (10 points) Prove that the total energy of the system is conserved.

According to Newton's eqn:  $m\ddot{x}(t) = F(x(t))$

$$E = \frac{m\dot{x}^2}{2} + V(x) \quad \text{where } V(x) \text{ is a potential}$$

such that  $F = -V'$

$$\begin{aligned} \frac{d}{dt} E &= m\ddot{x}\dot{x} + \frac{dV}{dt} = m\ddot{x}\dot{x} + V'\dot{x} \\ &= \dot{x}(m\ddot{x} + V') = \dot{x}(m\ddot{x} - F) = 0 \end{aligned}$$

∴  $E = \text{constant}$ .

Now, let  $F$  be given by  $F(x) := ae^{-x/\ell}$ , where  $a$  and  $\ell$  are nonzero real constants.

- 1.b (10 points) Give the Euler-Lagrange equation for the system, simplify it as much as possible, and show that it is equivalent to Newton's equation for this system.

$$\mathcal{L} = \frac{m\dot{x}^2}{2} - V(x)$$

$$\frac{\partial \mathcal{L}}{\partial x} = -V' \quad , \quad \frac{\partial \mathcal{L}}{\partial \dot{x}} = m\ddot{x}$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}} - \frac{\partial \mathcal{L}}{\partial x} = 0 \Rightarrow m\ddot{x} - (-V') = 0$$

$$\Rightarrow m\ddot{x} = -V' = F(x) = a\ell e^{-\frac{x}{\ell}}$$

$$\text{For } F = a\ell e^{-\frac{x}{\ell}} \quad V = - \int F dx = - \left( a\ell \frac{e^{-\frac{x}{\ell}}}{-\frac{1}{\ell}} + C \right) = a\ell e^{-\frac{x}{\ell}} - C$$

We can set  $C=0$  without loss of generality, so that  $V(x) = a\ell e^{-\frac{x}{\ell}}$  const.

- 1.c (10 points) Give Hamilton's equations for the system, simplify them as much as possible, and show that they are equivalent to the Newton's equation.

$$P = \frac{\partial \mathcal{L}}{\partial \dot{x}} = m\dot{x} \Rightarrow \dot{x} = \frac{P}{m}$$

$$H = \dot{x}P - \mathcal{L} = \frac{P^2}{m} - \left[ \frac{1}{2m} \left( \frac{P}{m} \right)^2 - V(x) \right]$$

$$= \frac{P^2}{2m} + V(x) \quad V(x) = a\ell e^{-\frac{x}{\ell}}$$

$$\begin{cases} \dot{x} = \frac{\partial H}{\partial P} = \frac{P}{m} \Rightarrow P = m\dot{x} \\ \dot{P} = -\frac{\partial H}{\partial x} = -V' \end{cases} \quad \begin{array}{c} \downarrow \\ \frac{d}{dt}(m\dot{x}) = -V' \\ \downarrow \qquad \downarrow \\ m\ddot{x} = F(x) \end{array}$$

$$\Rightarrow m\ddot{x} = a\ell e^{-\frac{x}{\ell}}$$

**Problem 2** (20 points) In Problem 1, take  $x(0) = 0$  and  $\dot{x}(0) = 0$ , and use conservation of total energy to obtain the position of the particle for  $t > 0$ . To get full credit you must simplify the expression you find for  $x(t)$  as much as possible.

Hint: For any real number  $b$ ,

$$\int \frac{dx}{\sqrt{b^2 - e^{-x/t}}} = \left(\frac{t}{b}\right) \ln \left( \frac{1 + \sqrt{1 - b^{-2}e^{-x/t}}}{1 - \sqrt{1 - b^{-2}e^{-x/t}}} \right) + c.$$

where  $c$  is an integration constant.

$$E = \frac{m \dot{x}^2}{2} + V(x) = \frac{m \dot{x}^2}{2} + \alpha \ell e^{-\frac{x}{\ell}}$$

$$\Rightarrow \dot{x} = \pm \sqrt{\frac{2}{m} (E - \alpha \ell e^{-\frac{x}{\ell}})}$$

$$\text{Because } \dot{x}(0) = 0 \Rightarrow x(0) = 0 \Rightarrow E = \alpha \ell \Rightarrow$$

$$\dot{x} = \pm \sqrt{\frac{2\alpha \ell}{m}} \sqrt{1 - e^{-\frac{x}{\ell}}} \quad \text{for } \dot{x} \in \mathbb{R} \Rightarrow \alpha \ell > 0$$

$$\Rightarrow \int_0^x \frac{dx'}{\sqrt{1 - e^{-\frac{x'}{\ell}}}} = \pm \sqrt{\frac{2\alpha \ell}{m}} (t - t_0) = \pm \sqrt{\frac{2\alpha \ell}{m}} t$$

$$\ell \ln \left( \frac{1 + \sqrt{1 - e^{-x/\ell}}}{1 - \sqrt{1 - e^{-x/\ell}}} \right) \Big|_0^x = \pm \sqrt{\frac{2\alpha \ell}{m}} t$$

$$\Rightarrow \ell \ln \left( \frac{1 + \sqrt{1 - e^{-x/\ell}}}{1 - \sqrt{1 - e^{-x/\ell}}} \right) = \pm \sqrt{\frac{2\alpha \ell}{m}} t$$

$$\ell > 0 \Rightarrow \pm = \text{sign } \ell = \text{sign } \alpha \Rightarrow \ln \left( \frac{1 + \sqrt{1 - e^{-x/\ell}}}{1 - \sqrt{1 - e^{-x/\ell}}} \right) = \frac{1}{\ell} \sqrt{\frac{2\alpha \ell}{m}} t$$

$$\Rightarrow \frac{1 + \sqrt{1 - e^{-x/\ell}}}{1 - \sqrt{1 - e^{-x/\ell}}} = \frac{e^{\beta t}}{1}$$

$$\Rightarrow \frac{1}{2} (1 + \sqrt{1 - e^{-x/\ell}}) = \frac{e^{\beta t}}{1 + e^{\beta t}} \Rightarrow 1 - e^{-\frac{x}{\ell}} = \left( \frac{e^{\beta t}}{1 + e^{\beta t}} - 1 \right)^2 = \left( \frac{e^{\beta t} - 1}{1 + e^{\beta t}} \right)^2$$

$$\Rightarrow x = -\ell \ln \left[ 1 - \left( \frac{e^{\beta t} - 1}{e^{\beta t} + 1} \right)^2 \right] = -\ell \ln \left[ \left( \frac{2}{e^{\beta t} + 1} \right) \left( \frac{2 e^{\beta t}}{e^{\beta t} + 1} \right) \right] = -\ell \left[ 2 \ln \left( \frac{e^{\beta t} + 1}{2} \right) - \beta t \right]$$

Let  $\beta := \sqrt{\frac{2\alpha}{m\ell}}$

**Problem 3** (20 points) Give the Hamilton-Jacobi equation for the system described in Problem 1. Show that the solution of this equation reproduces your response to Problem 2.

$$-\frac{\partial S}{\partial t} = H(x, S') = \frac{1}{2m} \left( \frac{\partial S}{\partial x} \right)^2 + V(x), \quad S' := \frac{\partial S}{\partial x}$$

$$\text{let } S = -tE + w(x) \Rightarrow p = \frac{\partial S}{\partial x} = w' := \frac{\partial w}{\partial x}$$

$$\Rightarrow E = \frac{w'^2}{2m} + \alpha l e^{-\frac{x}{l}}$$

$$\text{let } \tilde{x} := x(0), \quad \tilde{p} := p(0) \Rightarrow$$

$$E = \frac{\tilde{p}^2}{2m} + \alpha l e^{-\frac{\tilde{x}}{l}} \quad (8)$$

Also we have

$$w = \pm \int dx \sqrt{2m(E - \alpha l e^{-\frac{x}{l}})}$$

$$\tilde{p} = -\frac{\partial S}{\partial \tilde{x}} = -t \frac{\partial E}{\partial \tilde{x}} - \frac{\partial w}{\partial \tilde{x}} = \left( -t + \frac{\partial w}{\partial E} \right) \frac{\partial E}{\partial \tilde{x}}$$

$$\frac{\partial E}{\partial \tilde{x}} = -\alpha e^{-\frac{\tilde{x}}{l}} = -\alpha \neq 0$$

$$\tilde{x} = x(0) = 0$$

$$\tilde{p} = p(0) = 0$$

$$0 = \left( -t + \frac{\partial w}{\partial E} \right) \underbrace{\frac{\partial E}{\partial \tilde{x}}}_{-\alpha} \Rightarrow$$

$$\frac{\partial w}{\partial E} = -t$$

$$\pm \int dx \frac{m}{\sqrt{2m(E - \alpha l e^{-\frac{x}{l}})}} = -t$$

$$\Rightarrow \pm \frac{m}{\sqrt{2mE}} \int \frac{dx}{\sqrt{1 - \frac{\alpha l}{E} e^{-\frac{x}{l}}}} = -t$$

$$\text{From (8)} \quad E = \alpha l \quad \Rightarrow$$

$$\int \frac{dx}{\sqrt{1 - e^{-\frac{x}{l}}}} = \pm \sqrt{\frac{2\alpha l}{m}} t$$

This is precisely the eqn  
we derived in Problem 2.  
So it gives the same solution.

**Problem 4** Consider a particle moving on the real line whose motion is described by the Hamiltonian

$$H = \alpha e^{\mu t} q^2 p^{4/3} + \beta e^{-\mu t} p^{2/3},$$

where  $\alpha$ ,  $\beta$ , and  $\mu$  are positive real parameters, and  $q$  and  $p$  are respectively the position and momentum of the particle.

**4.a** (5 points) Find a function  $f(t)$  such that for every real number  $\nu$  the following coordinate transformation of the phase space of this particle is a canonical transformation:  $(q, p) \rightarrow (\tilde{q}, \tilde{p})$  where

$$\begin{aligned} l &= \left\{ \tilde{q}, \tilde{p} \right\} = \frac{\partial \tilde{q}}{\partial q} \frac{\partial \tilde{p}}{\partial p} - \frac{\partial \tilde{q}}{\partial p} \frac{\partial \tilde{p}}{\partial q} \\ &= - \left( \frac{1}{3} e^{\nu t} p^{-2/3} \right) \left( \frac{f(t)}{3} p^{2/3} \right) = - \frac{e^{\nu t} \frac{f(t)}{3}}{3} \\ &\Rightarrow \boxed{f(t) = -3 e^{-\nu t}} \end{aligned}$$

**4.b** (5 points) Choose  $\nu$  so that when you express the right-hand side of the above expression for  $H$  in terms of  $\tilde{q}$  and  $\tilde{p}$  it does not depend on  $t$  explicitly.

$$\begin{aligned} H &= \alpha e^{\nu t} \tilde{q}^2 \tilde{p}^{4/3} + \beta e^{-\nu t} \tilde{p}^{2/3} = \alpha e^{\nu t} \left( \frac{\tilde{p}}{f(t)} \right)^2 + \beta e^{-\nu t} \left( \frac{\tilde{q}}{e^{\nu t}} \right)^2 \\ &= \alpha e^{\nu t} \left( \frac{\tilde{p}^2}{9 e^{-2\nu t}} \right) + \beta e^{-\nu t} \left( \frac{\tilde{q}^2}{e^{2\nu t}} \right) \end{aligned}$$

$$\text{So take } \nu = -2\omega \Rightarrow \boxed{\omega = -\frac{\mu}{2}}$$

$$\Rightarrow H = \frac{\alpha \tilde{p}^2}{9} + \beta \tilde{q}^2$$

4.c (15 points) For the function  $f(t)$  and the value of  $\nu$  you find in Problems 4.a and 4.b.  $(q, p) \rightarrow (\tilde{q}, \tilde{p})$  is a canonical transformation. Find a generator of the form  $F(q, \tilde{q})$  for this canonical transformation.

Hint: Recall that

$$p = \frac{\partial F}{\partial q}, \quad \tilde{p} = -\frac{\partial F}{\partial \tilde{q}}, \quad \tilde{H} = H + \frac{\partial F}{\partial t}.$$

$$\begin{aligned} & \Downarrow \\ F &= \int p dq + g(\tilde{q}) \\ p &= (e^{-\omega t} \tilde{q})^3 = (e^{\frac{3\mu}{2}t} \tilde{q})^3 \\ &= e^{\frac{3\mu}{2}t} \tilde{q}^3 \end{aligned}$$

$$\Rightarrow F = \int e^{\frac{3\mu}{2}t} \tilde{q}^3 dq + g(\tilde{q}) = e^{\frac{3\mu}{2}t} \tilde{q}^3 q + g(\tilde{q})$$

$$\tilde{p} = -\frac{\partial F}{\partial \tilde{q}} \Rightarrow f(t) q | p^{2/3} = -3 e^{\frac{3\mu}{2}t} \tilde{q}^2 - g'(\tilde{q})$$

$$\begin{aligned} -3 e^{-\omega t} &\Downarrow \\ -3 e^{\frac{3\mu}{2}t} &= (e^{\frac{3\mu}{2}t} \tilde{q}^3)^{2/3} = e^{\mu t} \tilde{q}^2 \\ -3 e^{\frac{\mu}{2}t} & \end{aligned}$$

$$\Rightarrow -3 e^{\frac{3\mu}{2}t} \tilde{q}^2 = -3 e^{\frac{3\mu}{2}t} \tilde{q}^2 - g'(\tilde{q})$$

$$\Rightarrow g'(\tilde{q}) = 0 \Rightarrow g(\tilde{q}) = \text{const} =: g_0$$

$$\boxed{F(q, \tilde{q}) = e^{\frac{3\mu}{2}t} \tilde{q}^3 q + g_0}$$

4.d (5 points) Determine the transformed Hamiltonian  $\tilde{H}$  and simplify it as much as possible.

$$\tilde{H} = H + \frac{\partial F}{\partial t}$$

$$\frac{\partial F}{\partial t} = \frac{3}{2} \mu e^{\frac{3\mu t}{2}} \tilde{q} \tilde{q}^3$$

$$\tilde{q} = \frac{\tilde{p}}{f(t)} e^{-\frac{3\mu t}{2}} = \left( \frac{\tilde{p}}{-3e^{\frac{3\mu t}{2}}} \right) \left( -\frac{\tilde{q}}{e^{-\frac{3\mu t}{2}}} \right)^{-2}$$

$$= \frac{\tilde{p}}{-3 e^{\frac{3\mu t}{2}} \tilde{q}^2}$$

$$\Rightarrow \frac{\partial F}{\partial t} = -\frac{1}{2} \mu \tilde{q} \tilde{p}$$

$$\Rightarrow \boxed{\tilde{H} = \frac{\tilde{p}^2}{2} + \beta \tilde{q}^2 - \frac{1}{2} \mu \tilde{q} \tilde{p}}$$